

Molecular Dynamics Simulation Studies of CO₂ – [bmim][PF₆] Solutions: Effect of CO₂ Concentration

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DOI 10.1002/aic.11596

Published online August 27, 2008 in Wiley InterScience (www.interscience.wiley.com).

Molecular dynamics simulations have been carried out on CO₂ – [bmim][PF₆] mixtures using a refined atomistic potential model for the ionic liquid, at different concentrations of CO₂. The expansion in volume as a function of added CO₂ was found to agree well with experiments at all but the highest concentration. Significant concentration dependent differences in the radial distribution function of CO₂ around the anion have been observed. These differences have been attributed to the specific interaction between CO₂ and the anions. The diffusion coefficients of the ions and of CO₂ have been found to increase with increase in CO₂ concentration. The rotational relaxation of CO₂ molecule in solution is found to be biexponential, and the mean relaxation times decreases with increasing CO₂ concentration. © 2008 American Institute of Chemical Engineers AIChE J, 54: 2971–2978, 2008

Keywords: ionic liquids, supercritical carbon dioxide, molecular dynamics simulation, structure, dynamics

Introduction

Room temperature ionic liquids (RTILs), along with supercritical carbon dioxide (scCO₂), have been receiving considerable attention from the industry and academia recently.¹ scCO₂ is one of the well-studied systems through experimental^{2,3} and computational^{4–7} approaches. One of the interesting characteristics of scCO₂ is its liquid-like density and gas-like transport properties. RTILs, on the other hand, are polar and hence solvate a large number of organic and inorganic compounds. Some of the ILs also possess negligible vapor pressure at ambient conditions^{8,9} and a wide liquid range. With a large variety of cations and anions which can be used, in principle, an ionic liquid can be tailor made for a specific application.^{10,11} Possessing such interesting properties, these compounds have received a lot of attention in the recent

past.^{12–22} Supplementary to the experiments, a variety of computational studies have provided microscopic details of these systems.^{23–32}

scCO₂ has been used for extraction of products from RTILs.^{33–36} The interaction between the solute and the solvent in the mixture can impart new properties which the individual components may not possess.³⁷ Although CO₂ is remarkably soluble in imidazolium based ionic liquids,^{33,38–43} ionic liquids such as [bmim][PF₆] do not dissolve much in CO₂.^{33,38} The separation of organic compounds and water from ILs using CO₂ has been demonstrated.³⁵ There have been many efforts in the understanding of molecular level interactions between RTIL and CO₂.^{44–52}

By observing the effects of varying the anion among different ILs, Maginn and coworkers⁴⁶ have shown that the anion influences the solubility of CO₂ in ionic liquids to a greater extent compared with cation. Earlier, the nature of these anion–CO₂ interaction has been probed using vibrational spectroscopic methods.⁴⁴ Kazarian et al. have reported a weak Lewis acid–base interaction between CO₂ and the

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Table 1. Details of Simulations

CO ₂ (mol %)	Pressure (MPa)	Box Length (Å)	Density (kg m ⁻³)
0	0.0	44.30	1389
10	0.6	44.56	1388
30	2.0	45.38	1378
50	4.4	46.81	1360
70	15.0	49.92	1320

The uncertainty in the density is less than 0.1%.

anions of the ILs.⁴⁴ Computer simulation studies have demonstrated that CO₂ occupies preexisting void spaces in the IL that reorganize to accommodate the CO₂,⁴⁸ which rationalizes for the negligible increase in the molar volume of the mixture compared with pure ILs at low concentration of CO₂. X-ray diffraction experiments have been carried out by Kanakubo et al.⁴⁹ to examine the intermolecular structure of the solution and, in particular, the location and orientation of the CO₂ molecule with reference to the anion. Based on the analysis of experimental data, they have concluded that the CO₂ molecule may not be tangential to the PF₆ sphere as claimed in the computer simulation studies.^{46,48,53}

We have recently carried out ab initio MD simulations of the [bmim][PF₆] – CO₂ mixture at a CO₂ concentration of 70 mol % to study the intermolecular environment around CO₂ molecules in the solution.⁵³ In addition, these ab initio simulations enabled us to refine the interatomic potential model for [bmim][PF₆].⁵⁴ In the present study, we attempt to study the [bmim][PF₆] – CO₂ mixture at different concentrations of CO₂ using the refined potential model. We present the methods in the next section, followed by the results and end with conclusions drawn from our study.

Methodology and Simulation Details

Molecular dynamics simulations were carried out on [bmim][PF₆] – CO₂ mixtures with varying concentration of CO₂ at 300 K using the PINY_MD code.⁵⁵ The form of the interatomic potential and the potential parameters for the [bmim] and [PF₆] ions are adopted from our previous work.⁵⁴ The model has been found to predict the intermolecular structure in good agreement with the results of ab initio MD simulations.⁵³ It also predicts the diffusion coefficients of ions, density, and surface tension in excellent agreement with experimental data. For the CO₂ molecules, the EPM2 model⁵⁶ was used. Lorentz–Berthelot mixing rules were employed to obtain the Lennard–Jones parameters for cross-interactions. The model contains fully flexible [bmim] and [PF₆] ions. However, the bond distances in CO₂ were constrained in accordance with the EPM2 model.

Five different systems with 0, 10, 30, 50, and 70 mol % of CO₂ in 256 ion pairs of [bmim][PF₆] were simulated. The systems were first equilibrated under constant temperature and pressure (NPT) conditions, and later in the canonical (NVT) ensemble, with the volume determined as an average over the earlier NPT trajectory. The number of CO₂ molecules for the 10, 30, 50, and 70 mol % of CO₂ in solution turned out to be 28, 110, 256, and 595, respectively. The pressure of the system was chosen from experimental data.⁵⁷ In the canonical ensemble, the system was equilibrated for at

least 2 ns and a trajectory was generated for 2 ns in which the coordinates were stored every 0.9 ps for further analysis. The pressure applied in the isothermal–isobaric ensemble simulation, the converged box length for that pressure and the final densities are provided in Table 1. Other details are as provided earlier.⁵⁴

Nosé–Hoover chain thermostats⁵⁸ were used for temperature control with a thermostat time constant of 1 ps. The mean squared displacements of CO₂ molecules were found not to be sensitive to the use of thermostats during the analysis phase of the trajectory. Periodic boundary conditions consistent with cubic symmetry were employed to obtain the bulk behavior. Ewald summation method was used to handle Coulombic interactions with an α value of 0.339 Å⁻¹. Nonbonded interactions were calculated up to a distance cutoff of 13 Å and long range corrections were applied for the calculation of energy and pressure tensor. The multiple time step algorithm, r-RESPA⁵⁹ was used to integrate the equations of motion, in which the nonbonded interactions beyond 6 Å and within 13 Å were integrated with a time step of 3 fs and those within 6 Å with a time step of 1.5 fs. Torsional forces were computed every 0.75 fs and the bending and stretching degrees of freedom were integrated with a time step of 0.375 fs. The energy conservation was monitored through the simulations and was found to be three parts in 10⁵ over 1 ns.

Results and Discussions

Volume expansion

CO₂ is highly soluble in some of the ionic liquids. It is soluble up to 75 mol % in [bmim][PF₆].³³ Earlier studies⁴⁸ have indicated that CO₂ occupies preexisting void spaces in the ionic liquids, which is the reason for the low volume expansion on its addition to the RTILs. Voronoi analysis was used to conclude that the cation and the anion rearrange in the ionic liquid to form small number of large voids from large number of small voids to accommodate the added CO₂. It is thus important to verify if the current simulations predict the volume of the solution as a function of CO₂ concentration. The volume expansion predicted by the current model at different concentrations of CO₂ along with the experimental values⁵⁷ are provided in Table 2. The latter are the best fit values obtained by fitting the experimental data to a quadratic equation.

The experimental data were available only in the range of CO₂ concentrations between 0 and 56.7 mol %. The values obtained from our simulation are quite close to the experimental data within this domain. The volume expansion for

Table 2. Comparison of Volume Expansion of [bmim][PF₆] – CO₂ Mixture at Different Concentrations of CO₂ Obtained from Simulations and Experiment⁵⁷

Concentration of CO ₂ (mol %)	Volume Expansion (%) Computation	Volume Expansion (%) Experiment
10	1.8	1.5
30	7.5	6.9
50	18.0	16.8
70	43.1	31.2

The uncertainty in the computed value is less than 0.1%.

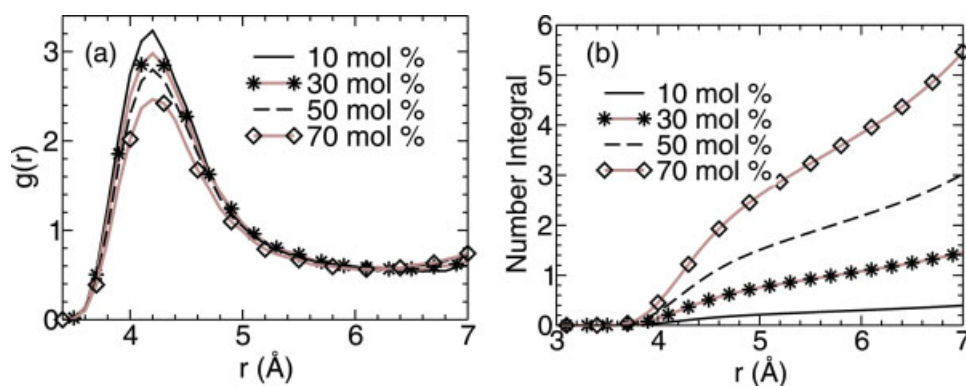


Figure 1. (a) Radial distribution functions of CO₂ around the PF₆ anion for different concentrations of CO₂ in the solution. Symbols are shown infrequently for clarity; (b) corresponding number integrals.

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the 70 mol % mixture is quite high in the simulations than the best fit to the experimental data. One of the reasons for this difference may be due to errors in the extrapolation procedure adopted for obtaining the experimental data.

Radial distribution functions

PF₆ – CO₂ RDF. The anion–CO₂ radial distribution functions (RDF) and the corresponding running number integrals are presented in Figure 1. The RDF peaks at 4.2 Å at all the concentrations studied. This peak position matches exactly with that reported by Huang et al.⁴⁸ The first minimum is present at 6.4 Å for the 50 and 70 mol % mixtures, whereas it is present at 6.2 Å for the 10 and 30 mol % mixtures. Figure 1 clearly shows that there is no shift in the peak positions with the change in the CO₂ concentration, although the peak heights are different. The coordination numbers up to 6.3 Å (up to the location of first minimum) are 0.3, 1.2, 2.4, and 4.2, respectively, for the 10, 30, 50, and 70 mol % mixtures. The increase in the coordination number with increase in the CO₂ concentration is reasonable. The coordination number up to the first minimum suggests that the CO₂ molecules interact with at least two of the anions. This is evident as the number of CO₂ that are within the first

coordination shell is almost thrice that of the average number of CO₂ per anion available in the system. The coordination shell of an anion is not exclusive, that is, the first coordination shell of one anion can overlap with that of the other. In the case of 10, 30, 50, and 70 mol % mixtures, the coordination numbers are found to be 3, 3, 2.5, and 2 times greater than the average number of CO₂ per anion, respectively. These suggest that the number of CO₂ that can be present in the first coordination shell around the anion is limited by CO₂–CO₂ repulsion. It has been demonstrated in our previous studies that CO₂ occupies the octahedral voids around the anions.⁵³ Although there are eight such voids, the number of CO₂ present around the anion is less than eight, suggesting that steric repulsion between the carbon dioxide molecules or the cation–anion intermolecular structure influence the CO₂ coordination at higher concentrations of CO₂.

PF₆ – PF₆ RDF. Figure 2 presents the PF₆ – PF₆ RDF and the number integrals. The anion–anion radial distribution function varies significantly with increase in CO₂ concentration. Slight differences can be observed on addition of even 10 mol % of CO₂ to the ionic liquid. The differences also increase with CO₂ concentration. The position of the first peak in the g(r) shifts from 6.5 Å in the case of pure [bmim][PF₆] to a value of 8.1 Å at the highest concentration

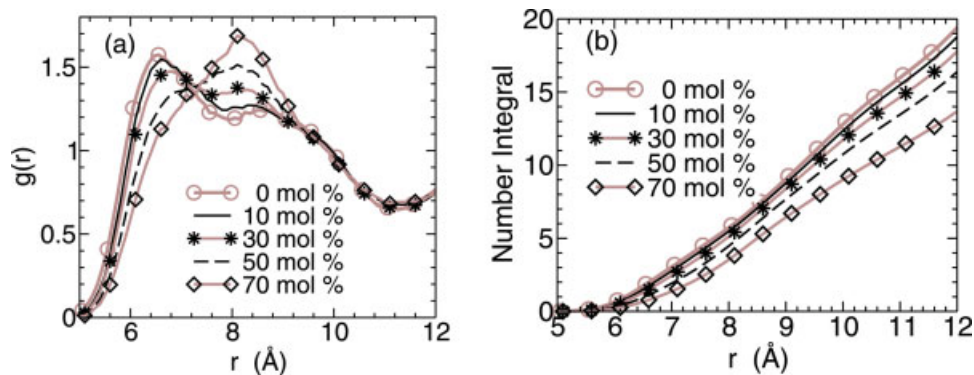


Figure 2. (a) PF₆–PF₆ radial distribution functions at various CO₂ concentrations in the mixture; (b) corresponding number integrals.

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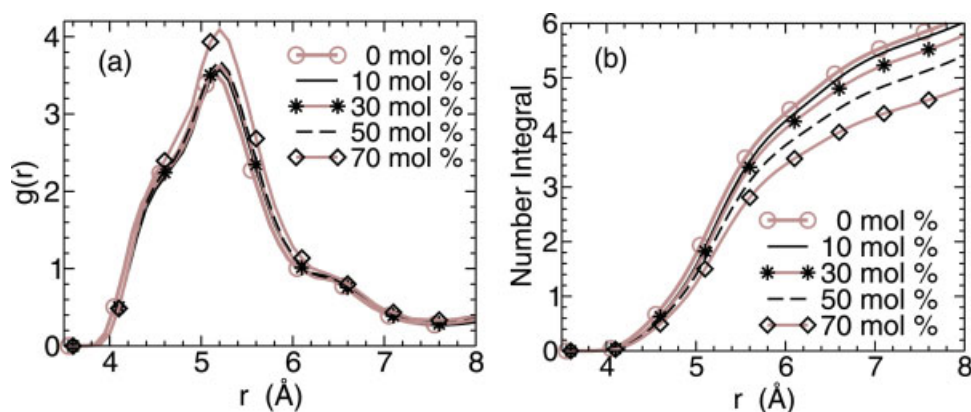


Figure 3. (a) Radial distribution functions of PF₆ anions around the cation ring center at various concentrations of CO₂ in the mixture; (b) corresponding number integrals.

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of CO₂. The latter is comparable with the value of 7.95 Å, which was observed by Huang et al.⁴⁸ as the position of the first peak. The first peak broadens too with increasing concentration of CO₂. Anions move away from each other with increase in CO₂ concentration in the mixture. CO₂ molecules are preferentially located around the anions, and hence with an increase in their concentration, anions tend to be separated further. However, the position of the first minimum in the $g(r)$ located at 11.2 Å is unaltered with increase in the CO₂ mole fraction.

As the anions move away from each other with introduction of CO₂, it is expected that the anion–anion coordination number will be lesser in case of mixture with highest concentration of CO₂ (see Figure 2b). The coordination numbers up to 7.2 Å (location where the intensity of all the RDF peaks is the same) are 3.3, 2.9, 2.3, and 1.7, respectively, for the 10, 30, 50, and 70 mol % mixtures. The same quantities up to the first minimum (at 11.2 Å) are 16.0, 15.2, 14.0, and 11.7, respectively.

[bmim] – PF₆ RDF. The cation–anion RDF is presented in Figure 3a, where the geometric center of the imidazolium ring is taken as the position of the cation. The corresponding

number integrals are shown in Figure 3b. RDFs at all concentrations of CO₂ are similar. All of them peak at 5.6 Å with their first minima at 7.6 Å. The coordination numbers up to the first minimum are 5.7, 5.5, 5.1, and 4.6, respectively, for the 10, 30, 50, and 70 mol % mixtures of CO₂ – [bmim][PF₆]. The coordination number decreases with the increase in CO₂ concentration. This is similar to that seen in anion–anion RDF, where the specific interaction between anion and CO₂ was said to be the reason for the decrease in coordination number. However, it should be noted that the decrease in coordination number in the case of cation–anion is smaller when compared with that for the anion–anion. This suggests that the anion interaction with CO₂ is mainly responsible for the separation of ions.

The large changes observed in the behavior of anion–anion RDFs and the lack thereof in the anion–cation RDFs, as a function of CO₂ concentration provides evidence for the interaction between CO₂ and anion.

CO₂–Ion RDF. We shall now examine the CO₂–anion interaction from the perspective of CO₂. In Figure 4a, we show the RDF of CO₂–Phosphorus and that of CO₂–C_r at two different concentrations of CO₂ in the mixture. C_r is the

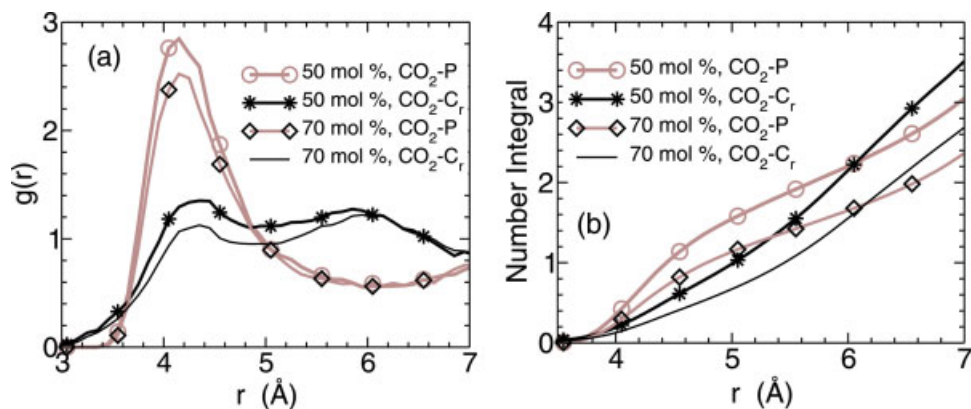


Figure 4. (a) RDFs of C(CO₂)–P and C(CO₂)–C_r at various concentrations of CO₂ in the solution and (b) corresponding number integrals.

C_r is the carbon atom present in between the two nitrogen atoms of the imidazolium (cation) ring. [Color figure can be viewed in the online issue, which is available at www.interscience.wiley.com.]

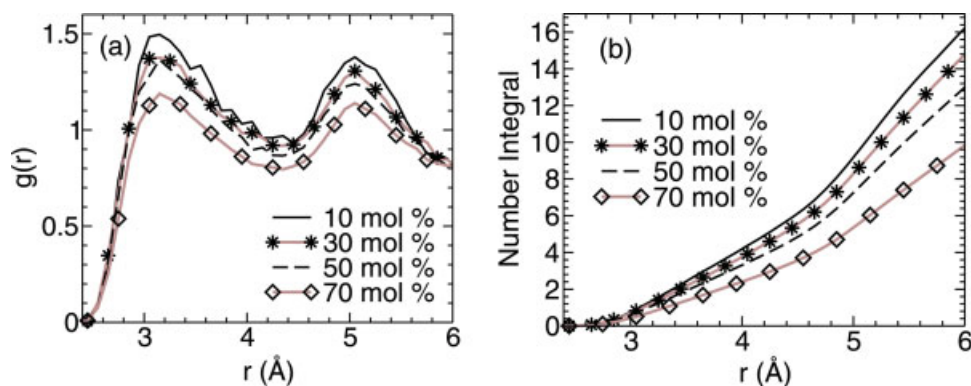


Figure 5. (a) C(CO₂)-fluorine radial distribution functions and (b) corresponding number integrals.

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ring carbon present in between the two ring nitrogens of the cation. Our ab initio MD simulations⁵³ had shown the presence of a weak hydrogen bond between the acidic hydrogen (which is covalently bonded to C_r) and the oxygen of CO₂. Hence, it is pertinent that we compare the CO₂-P RDF against the CO₂-C_r one. From the magnitudes of the first peak, it is obvious that carbon dioxide has a larger preference for the anion than the cation (here, the C_r site). This conclusion is confirmed from the number integrals shown in Figure 4b. In the 50 mol % mixture, at a distance of 5 Å, the number of anions surrounding a carbon dioxide is 1.6, whereas the number of cations (C_r site) is only 1.0.

The specific nature of interaction between carbon dioxide and the anion is of Lewis acid–base type. In fact, this is true not only of the PF₆ anion studied here, but for a variety of CO₂-anion complexes.⁶⁰ The carbon atom of CO₂ (Lewis acid) interacts favorably with the lone pair on the fluorine of the anion (Lewis base). The C(CO₂)-F RDF and the corresponding number integrals are shown in Figure 5 for all the mixtures studied. The strong peak in the RDFs present at 3.1 Å signifies specific interaction. At 10 mol % of CO₂, the number of fluorine atoms surrounding C(CO₂) within a distance of 4.2 Å (the first minimum in the RDF) is around 5. This implies that the first coordination shell of a CO₂ molecule contains fluorine atoms from more than one PF₆ anion. Thus, CO₂ molecules mediate (and possibly diminish) ion–ion interactions.

Diffusion coefficients

As an example, the mean squared displacement of anion and cation in the 30 mol % mixture that is obtained as an average over 5 independent configurations is presented in Figure 6, along with error bars. It can be seen from the figure that CO₂ molecules diffuse reasonably well. However, the ions have not diffused much within the time scale of the MSD plot. Experiments have shown a remarkable decrease in the viscosity of RTIL-CO₂ solution as a function of increasing CO₂ concentration.^{38,45,61} Maginn and coworkers^{62,63} as well as the group of Voth¹⁸ have demonstrated the slow dynamics of ions in RTILs. If one plots the MSD vs. time data on a double log graph, the slope of the curve (denoted by β) should be unity for diffusive behavior. In the 30 mol % mixture, the values of β at 500 ps were found to be 0.8, 0.7, and

0.98 for the cation, anion, and CO₂, respectively, whereas in the 70 mol % solution, the values were all above 0.95. Diffusion coefficients were calculated as the long-time regime slope of mean squared displacement curve (between 400 and 500 ps), for all concentrations of CO₂. In view of the fact that β is slightly less than unity, the diffusion coefficient values should be construed as apparent estimates. The exact values are likely to be slightly higher.⁶²

The diffusion coefficients of the ions as well as that of CO₂ molecules are found to increase with CO₂ concentration. It can be seen from the figure that the cation diffuses faster than the anion in [bmim][PF₆] – CO₂ mixtures which is also observed in experiments⁶⁴ as well as in simulations^{30,54} of pure ionic liquids. The values of diffusion coefficients at various concentrations of CO₂ are presented in Table 3. In the pure IL, the apparent diffusion coefficients obtained from our model agrees within 20% of the experimental value.⁵⁴ Experimental data for diffusion coefficients in CO₂-[bmim][PF₆] solutions are not available. However, as mentioned earlier, the increased diffusivity of ions and of CO₂ in these solutions

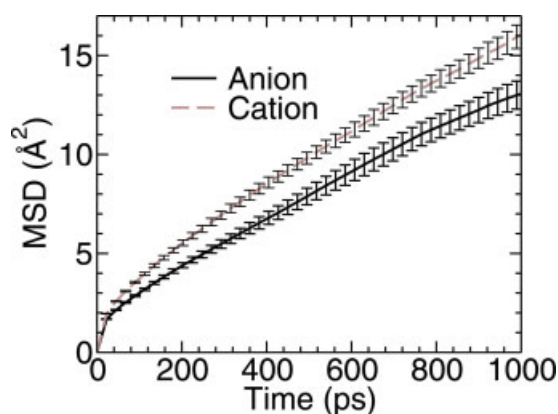


Figure 6. Mean squared displacement of anion and cation at 300 K in the 30 mol % mixture of CO₂-[bmim][PF₆] obtained as an average over five independent blocks, shown along with error bars.

[Color figure can be viewed in the online issue, which is available at www.interscience.wiley.com.]

Table 3. Apparent Diffusion Coefficients of Anion, Cation, and CO₂ in the [bmim][PF₆] – CO₂ Mixture Obtained from Simulations at Different Concentration of CO₂

CO ₂ (mol %)	D_- ($\times 10^{-12} \text{ m}^2 \text{ s}^{-1}$)	D_+ ($\times 10^{-12} \text{ m}^2 \text{ s}^{-1}$)	D_{CO_2} ($\times 10^{-12} \text{ m}^2 \text{ s}^{-1}$)
0	4.7	6.7	–
10	8.8	12.4	201.0
30	19.1	22.5	386.9
50	54.8	64.3	678.5
70	146.2	158.3	1453.8

The estimated uncertainty in the diffusion coefficient is 10%.

observed by us is consistent with the experimentally observed decrease in the viscosity of these solutions.^{38,45,61}

The variation of diffusion coefficients with concentration of CO₂ along with the best fit for anions and cations is presented in Figure 7. The increase in diffusion of ions is slow at low CO₂ concentrations but becomes steeper as one proceeds toward higher concentration of CO₂. The diffusion coefficients were fitted to a cubic equation of the form,

$$D_{\pm} = a_0 + a_1x + a_2x^2 + a_3x^3 \quad (1)$$

where a_0 , a_1 , a_2 , a_3 are constants and x is the mole fraction of CO₂ in the solution. The best fit values to curves that can be used to obtain diffusion coefficients values of ions and CO₂ at intermediate concentrations of CO₂ in the mixture are provided in Table 4. Huang et al.⁴⁸ have carried out MD simulations of solutions of [bmim][PF₆] containing 76.6 mol % of CO₂. The diffusion coefficients of anion, cation, and CO₂ obtained using the cubic fit to our data at 76.6 mol % are 196, 210, and 1920 (all in units of $10^{-12} \text{ m}^2 \text{ s}^{-1}$), respectively, whereas the values reported by Huang et al. are 68, 77, and 400 (all in units of $10^{-12} \text{ m}^2 \text{ s}^{-1}$), respectively. In the absence of experimental values of diffusion coefficients in these mixtures, the differences could only be attributed to differences in the force fields.

Orientalional relaxation

The reorientational time correlation function for the OCO vector of CO₂ gives us an idea of the rotational relaxation

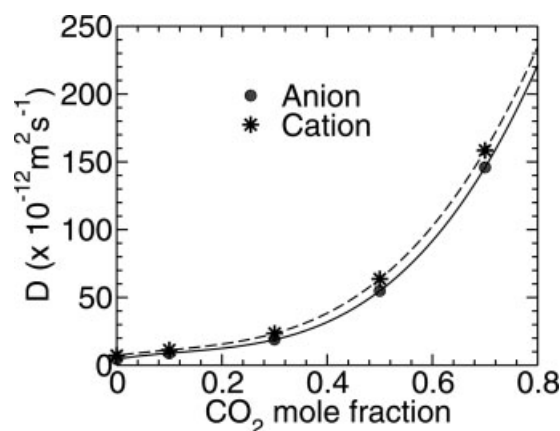


Figure 7. Variation of diffusion coefficient of ions with concentration of CO₂ in the solution.

The lines are cubic fits as expressed in Eq. 1.

Table 4. Values from the Best Fit to Eq. 1 for Diffusion Coefficients of Anion, Cation and CO₂

Ion/Molecule	a_0	a_1	a_2	a_3
Anion	4.6	59.5	–221.3	606.8
Cation	7.2	50.8	–155.5	559.6
CO ₂	2.4	2507.0	–6582.7	8521.2

times involved. Figure 8 show this quantity in the solution at different concentrations of CO₂. It can be seen that the backbone vector of CO₂ decorrelates faster in the case of mixtures with higher concentration of CO₂. The correlation was found to be lost completely before 120 ps in all the systems studied. These correlation functions show biexponential decay with two different relaxation times, one of the order of picoseconds and the other, in tens of picoseconds. The amplitude and the relaxation times obtained by fitting the data (up to 120 ps) to a constrained biexponential function of the form

$$y = a \exp(-t/\tau_1) + (1 - a) \exp(-t/\tau_2) \quad (2)$$

are tabulated in Table 5.

The faster relaxation for the systems with higher concentration of CO₂ can be observed at both timescales. The total relaxation is almost equally divided between these two time scales, which is evident from the amplitudes that are close to 0.5. The variation of the relaxation times with the CO₂ concentration shows a near linear behavior for both the timescales and also their amplitude weighted average. In summary, the CO₂ rotation is less hindered in the [bmim][PF₆] – CO₂ mixtures with higher mole fraction of CO₂, due possibly to reduced ion–ion interactions.

Conclusions

Molecular dynamics simulation studies have been carried out on [bmim][PF₆] – CO₂ mixtures at varying concentrations of CO₂. The expansion in the volume of the solutions

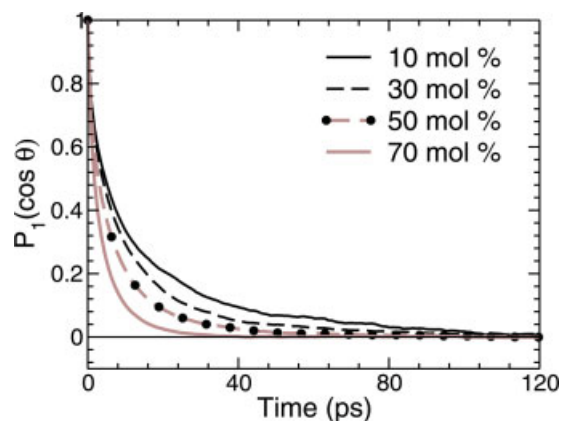


Figure 8. Reorientational time correlation function of the OCO vector of CO₂ at various concentrations of CO₂ in the solution.

[Color figure can be viewed in the online issue, which is available at www.interscience.wiley.com.]

Table 5. Amplitude of Relaxation and the Relaxation Times for the Reorientation of OCO Vector for Different Mole Fractions of CO₂ in [bmim][PF₆] – CO₂ Mixture

CO ₂ (mole %)	<i>a</i>	τ_1 (ps)	τ_2 (ps)
10	0.539	2.560	26.53
30	0.474	1.882	16.96
50	0.479	1.574	11.57
70	0.463	1.035	6.140

The statistical uncertainty in the time constants is 10%.

determined from simulations agree well with experimental data. Trends in the radial distribution of anions around an anion as well as that of anions around the cation indirectly suggest specific anion–CO₂ interaction. The nature of such an interaction (Lewis acid–base type) has been documented by us earlier.^{53,60} This aspect is also seen in the anion–CO₂ radial distribution function. The diffusion coefficient of the ions and of CO₂ increases with increase in CO₂ concentration. This nonlinear increase has been fitted to a cubic equation, which can be used to obtain the diffusion coefficients at intermediate values of CO₂ concentrations that have not been studied. The reorientational relaxation of the CO₂ molecule is found to be biexponential in time. The mean relaxation time decreases with increasing concentration of CO₂.

It is hoped that molecular simulations would considerably help in a microscopic understanding of gas solubilities in ionic liquids.⁶⁵

Acknowledgments

The research reported here is supported by the Department of Science and Technology, Government of India. B.L.B. thanks CSIR, India, for a research fellowship.

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Manuscript received Jan. 28, 2008, revision received Apr. 21, 2008, and final revision received Jun. 20, 2008.